

# Spin dynamics of water ice and the OPR of gaseous water desorbed from ice

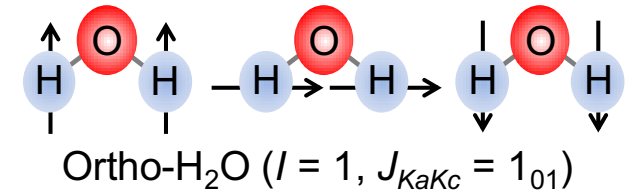
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Institute of Low Temperature Science, Hokkaido University



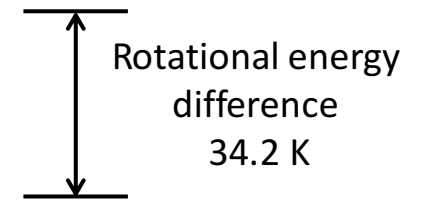
Nuclear Spin Effects in Astrochemistry, 2-4 May 2017,  
Grenoble, France

# Ortho and para water

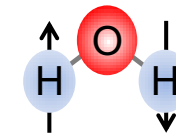
(1) Ortho-H<sub>2</sub>O and para-H<sub>2</sub>O must exist in different rotational states. (Proton is a fermion)



(2) Nuclear spin conversion is very slow in the gas phase by radiation or nonreactive collisions.

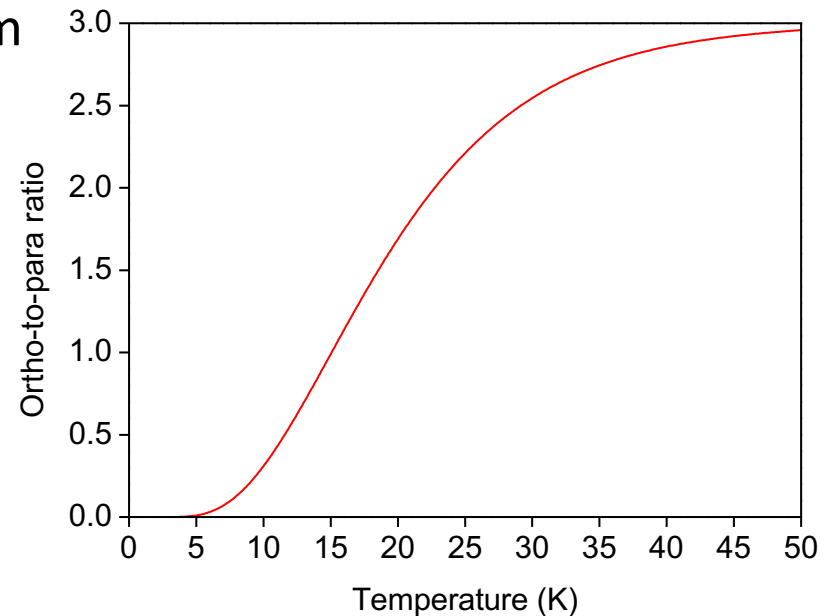


Miani and Tennyson, JCP (2004)., Cacciani et al., Phys. Rev. A (2012).



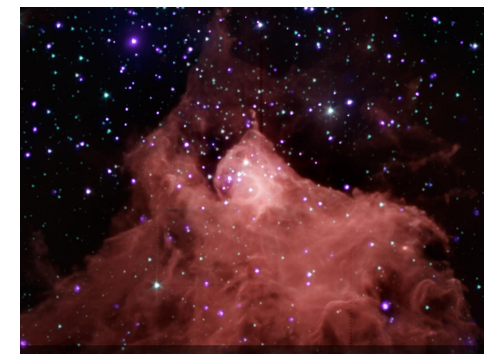
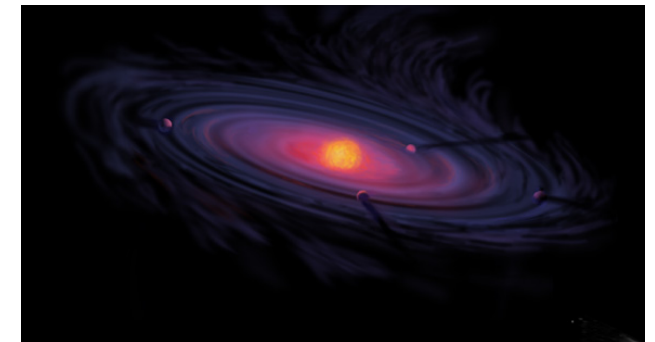
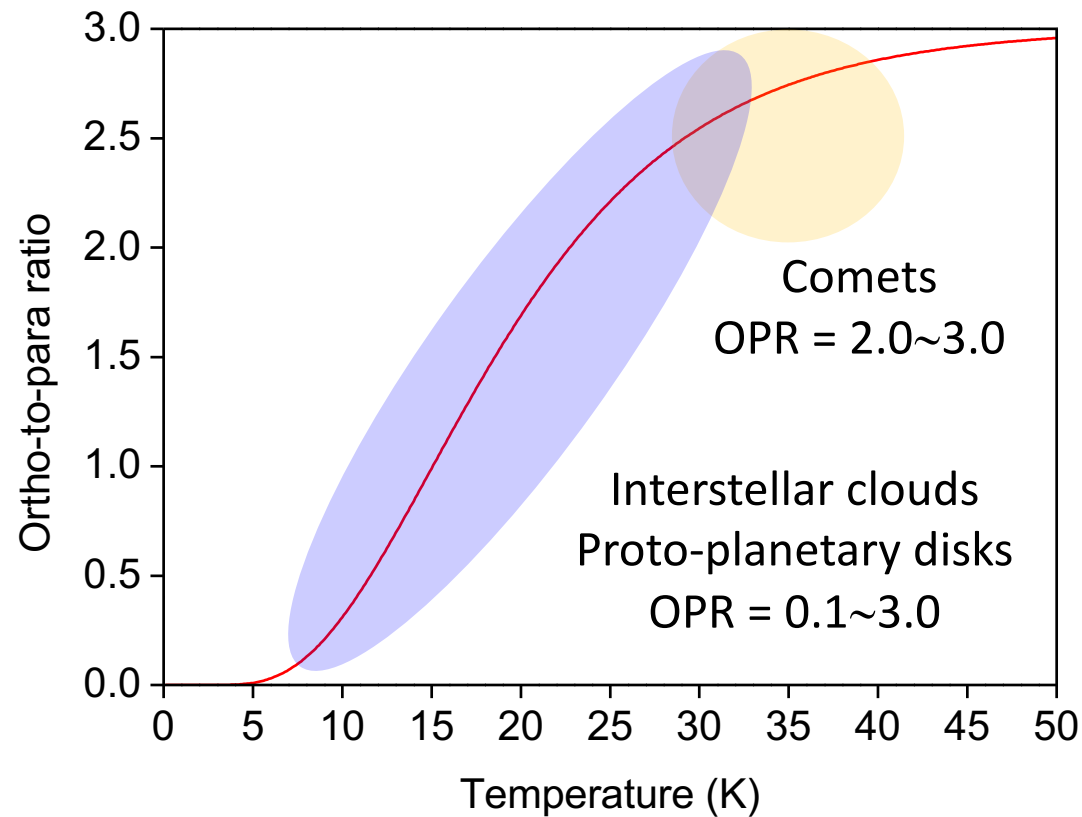
Spin temperature ( $T_{\text{spin}}$ ): Ortho-to-para ratio (OPR) assuming thermodynamic equilibrium

$$\text{OPR} = \frac{\boxed{\text{Spin and rotational degeneracy}} \sum (2J + 1) \exp \left[ \frac{\boxed{\text{Rotational energy}}}{k_B T_{\text{spin}}} \right]}{\sum (2J + 1) \exp \left[ \frac{-E_p(J_{Ka, Kc})}{k_B T_{\text{spin}}} \right]}$$



Because of  $\Delta E_{\text{rot}}$  ( $23.8\text{cm}^{-1}$ , 34.2 K), para-H<sub>2</sub>O ( $J_{Ka, Kc} = 0_{00}$ ) enrichment below 50 K.

The anomalous OPRs of gaseous H<sub>2</sub>O have been observed in space.



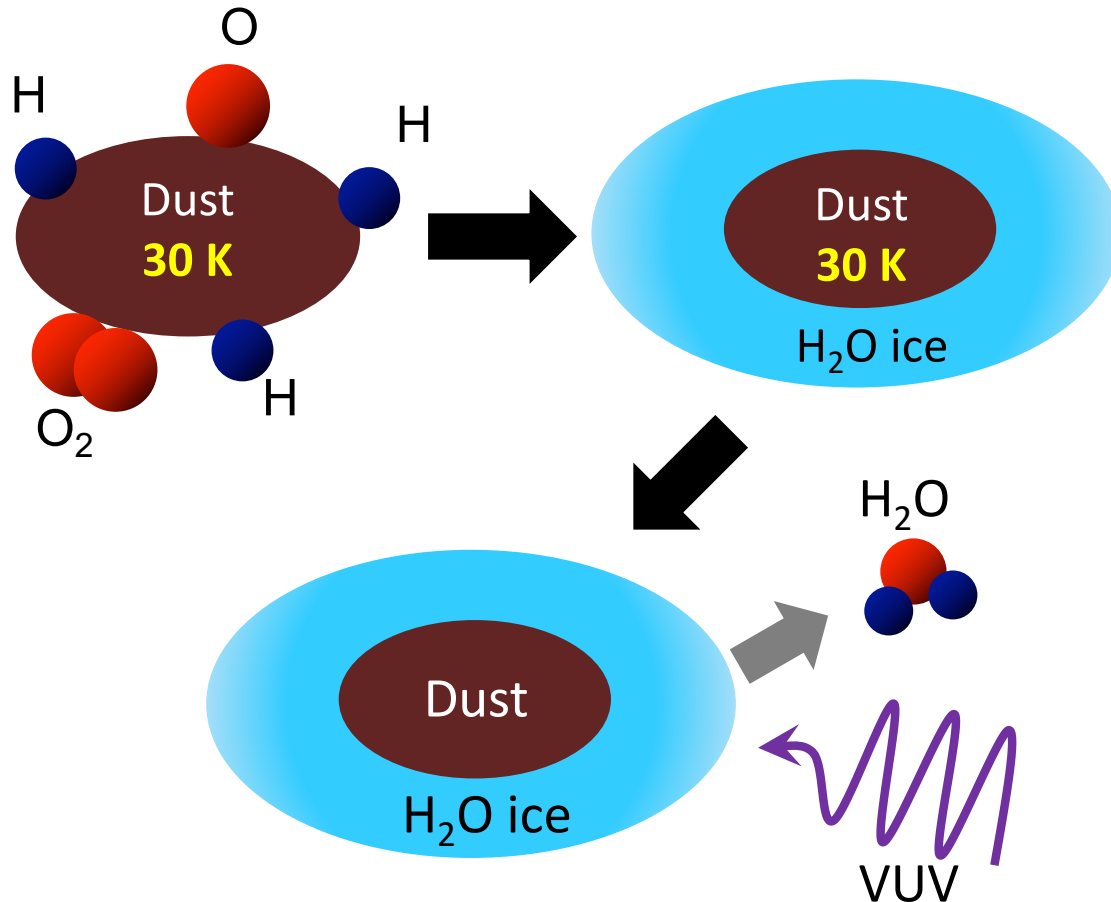
**The meaning of the observed  $T_{\text{spin}}$  remains a topic of continuing debate.**

Lis et al., *Astron. Astrophys.* 521, L26 (2010)., Hogerheijde et al., *Science* 334, 338 (2011).  
Choi et al., *Astron. Astrophys.* 572, L10 (2014)., Willacy et al., *Space Sci. Rev.* 197, 151 (2015).

# Hypothesis: $T_{\text{spin}}$ reflects the formation temperature of ice?

## Interstellar clouds / Proto-planetary disks

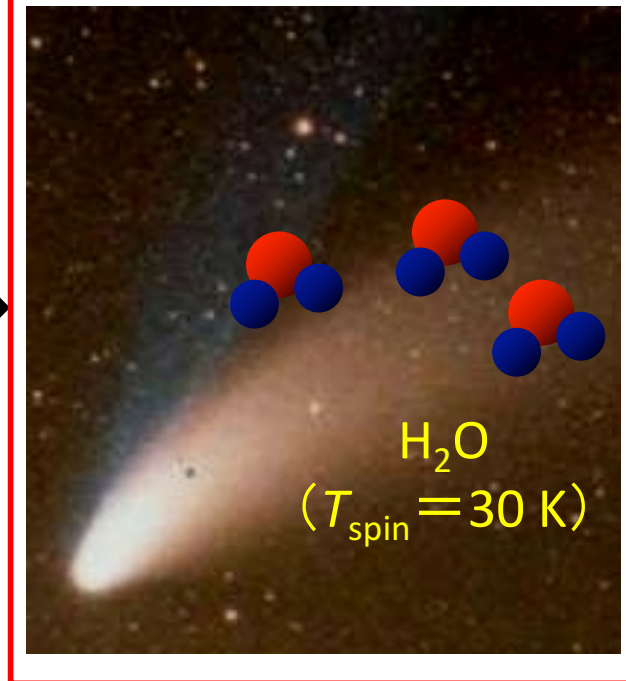
H<sub>2</sub>O ice is formed through dust surface reactions, e.g., H atom addition to O, O<sub>2</sub>, and OH.



## Photodesorption

Ice-covered dust (< 50 K) is too cold for the thermal desorption.

## Comets (Thermal desorption)



**What is the real meaning of the OPR?**

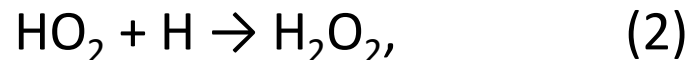
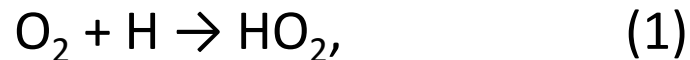
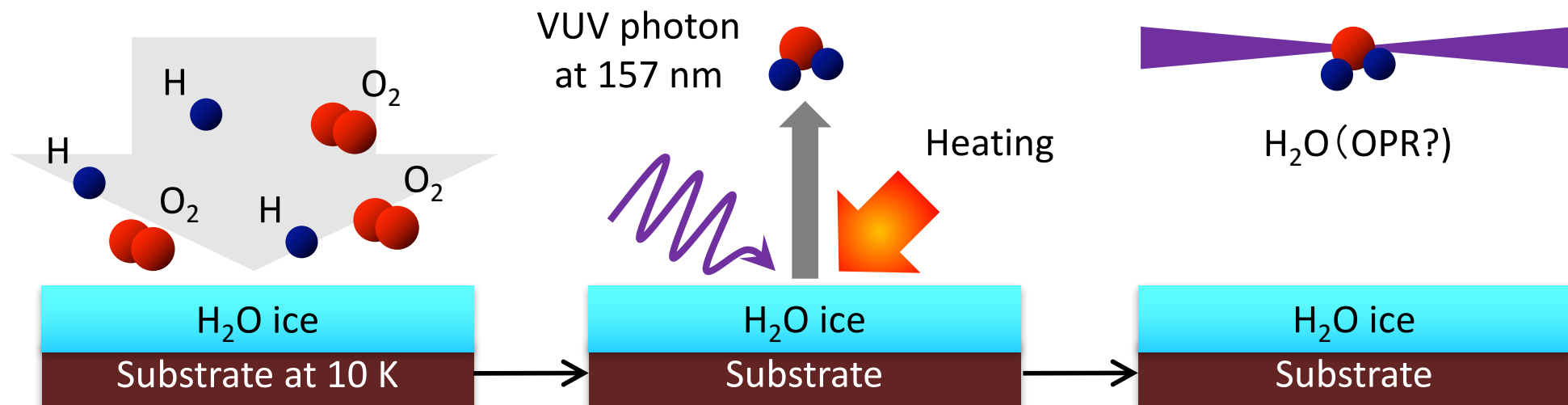


**Measurements of the OPR of H<sub>2</sub>O desorbed from ice.**

Test the relation between  $T_{\text{spin}}$  and the ice formation temperature:  
 $\text{H}_2\text{O}$  ice was produced in situ through the hydrogenation of solid  $\text{O}_2$  at 10 K.

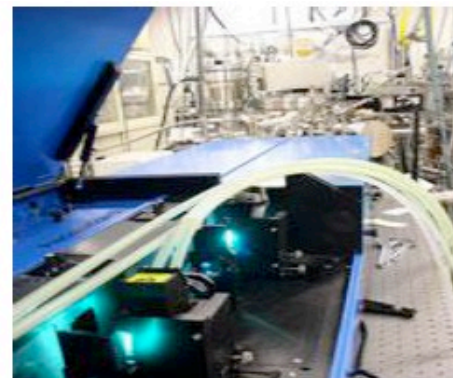
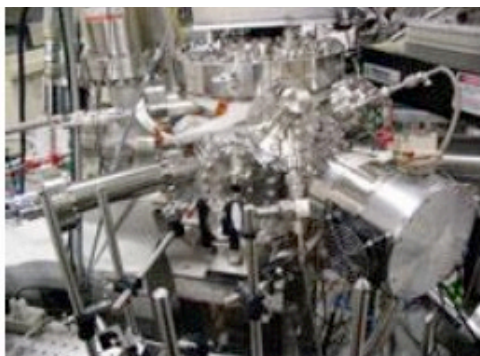
In situ production of  $\text{H}_2\text{O}$  ice at 10 K  
 by co-deposition of H and  $\text{O}_2$

Resonance-Enhanced MultiPhoton  
 Ionization (REMPI)

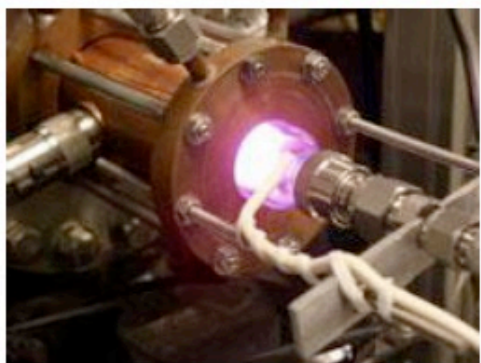


Surface reactions (3) – (5) are  
 the formation processes of  
 interstellar  $\text{H}_2\text{O}$  ice.

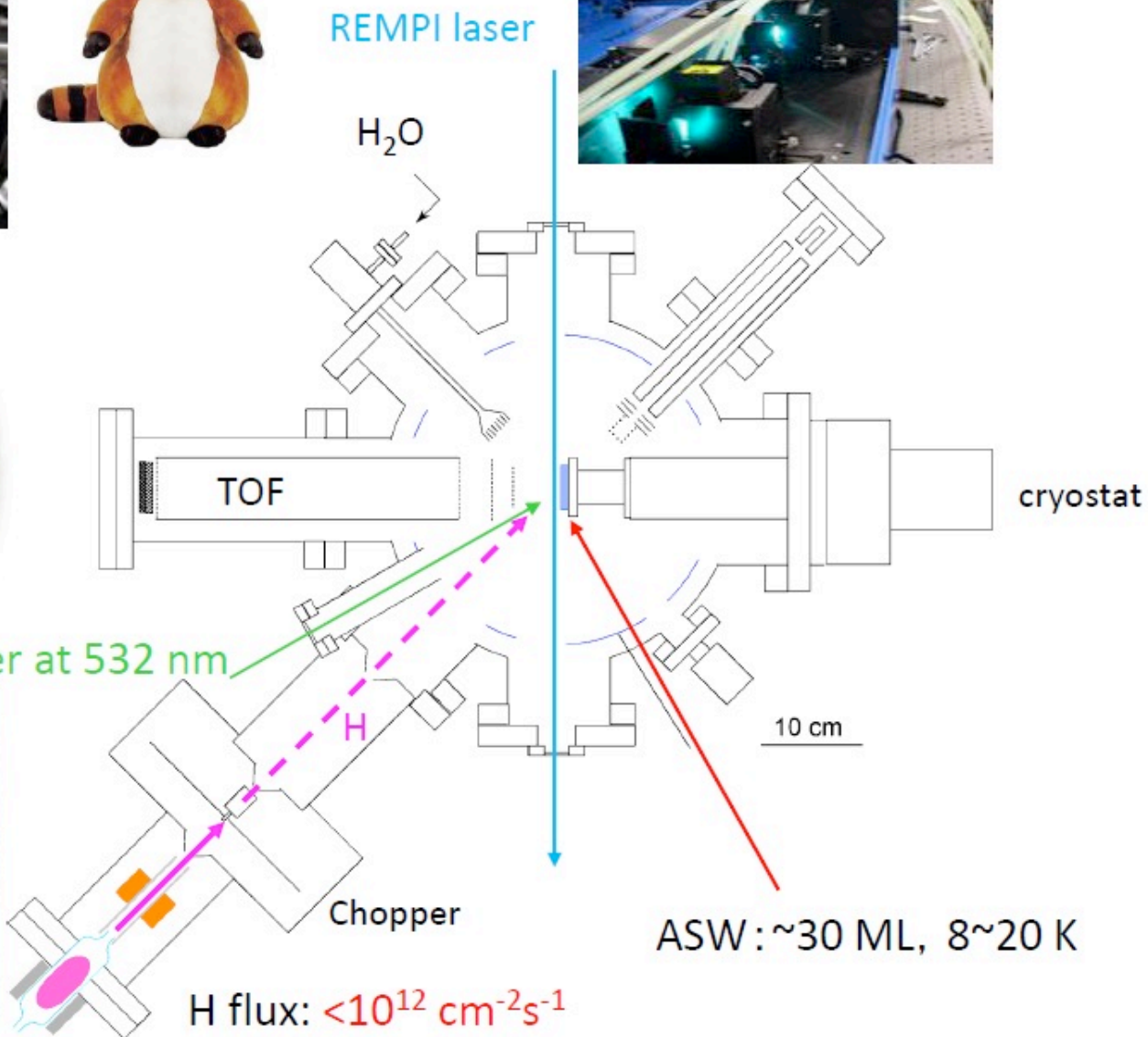
# Experiment with RASCAL



PSD laser at 532 nm



Microwave H or D atom source



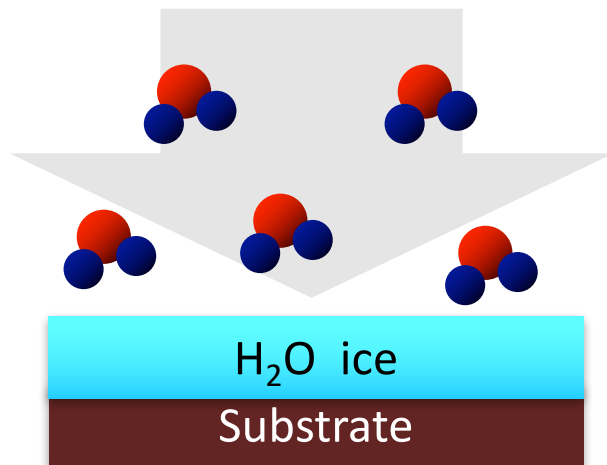


Infrared reflection-absorption spectra at 4000–800  $\text{cm}^{-1}$ .

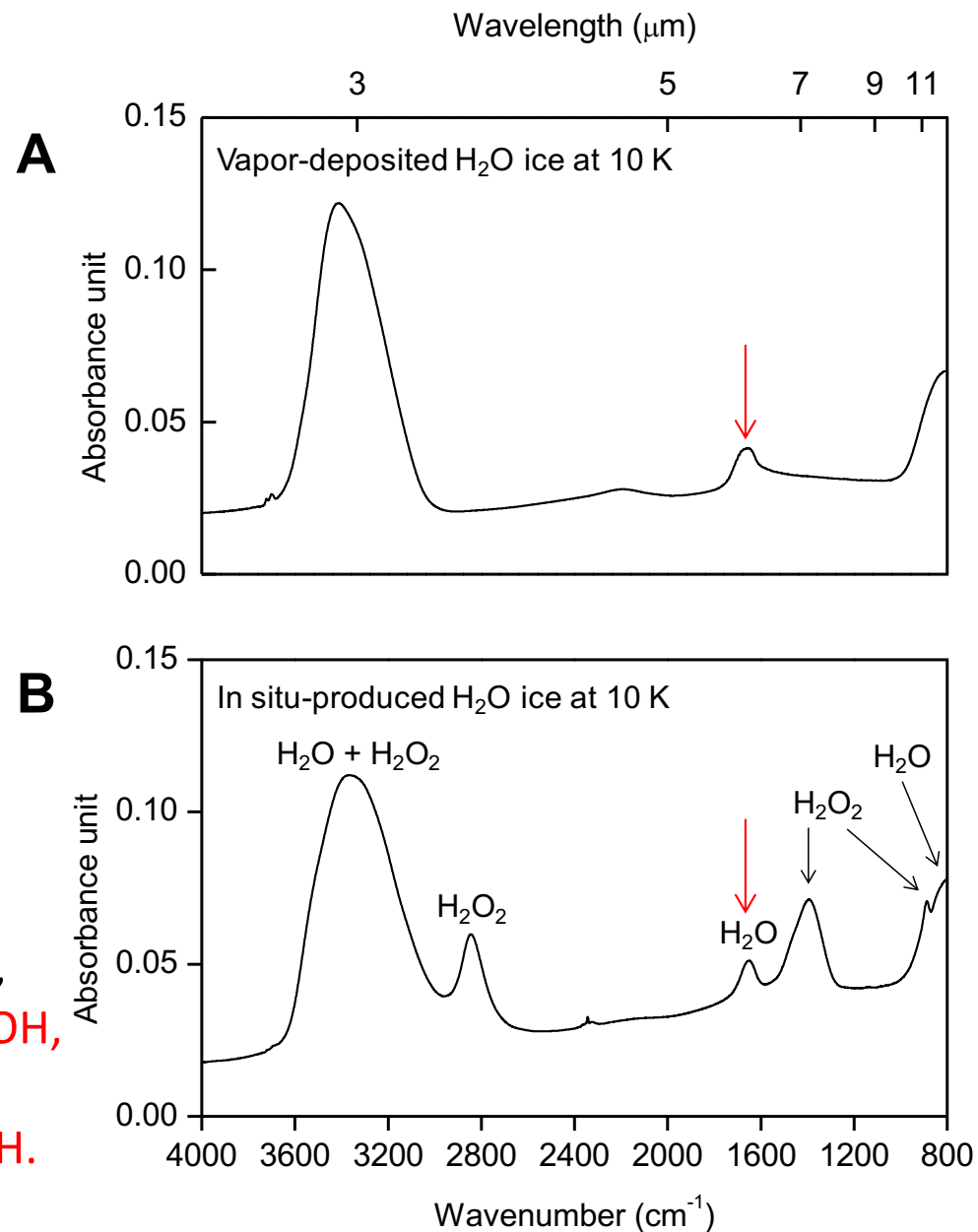
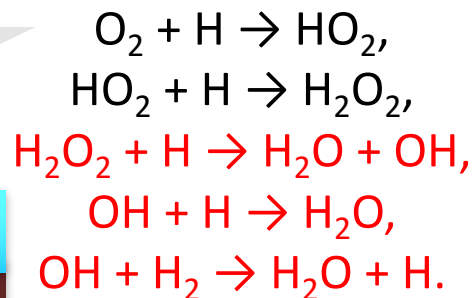
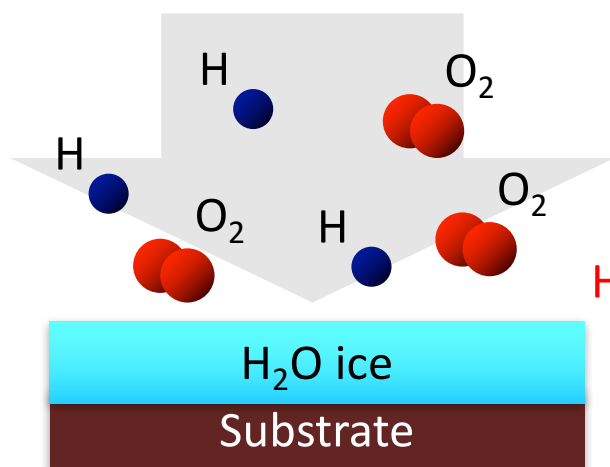
(A) Vapor-deposited  $\text{H}_2\text{O}$  ice at 10 K.

(B)  $\text{H}_2\text{O}$  ice produced by co-deposition of  $\text{O}_2$  with atomic H at 10 K for 420 min.

(A) Vapor-deposited  $\text{H}_2\text{O}$  ice at 10 K.



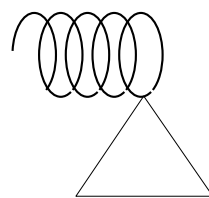
$\text{H}_2\text{O}$  ice produced by hydrogenation of  $\text{O}_2$  at 10 K



Detector

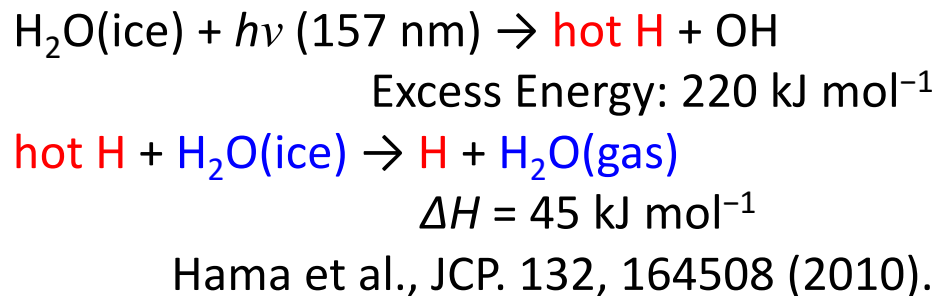
“kick-out” photodesorption: (TOF mass + MCP)

Momentum transfer from a hot H atom  
photodissociated from a neighboring H<sub>2</sub>O.

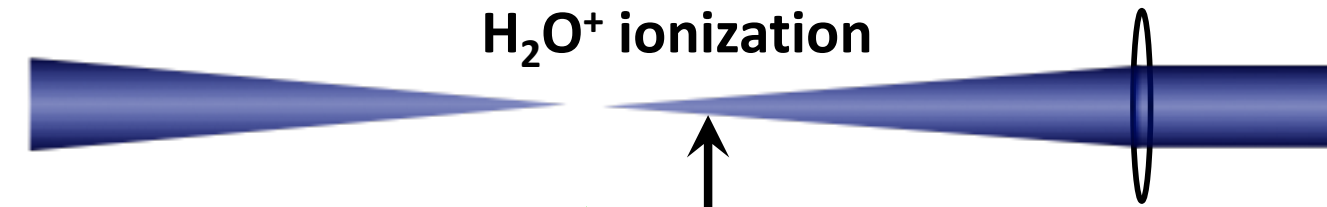


REMPI laser (248-249 nm)  
Focused (1 mJ pulse<sup>-1</sup>)

Ortho/Para state selective  
detection



H<sub>2</sub>O<sup>+</sup> ionization



157 nm

~1 mm

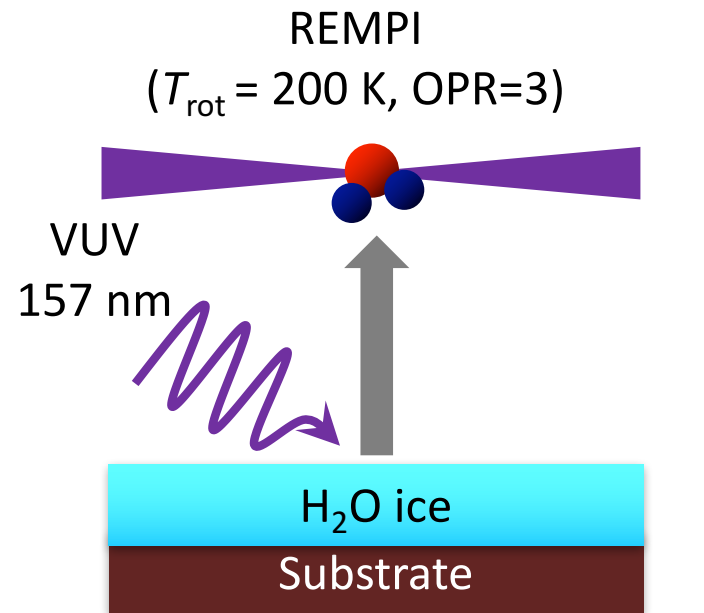
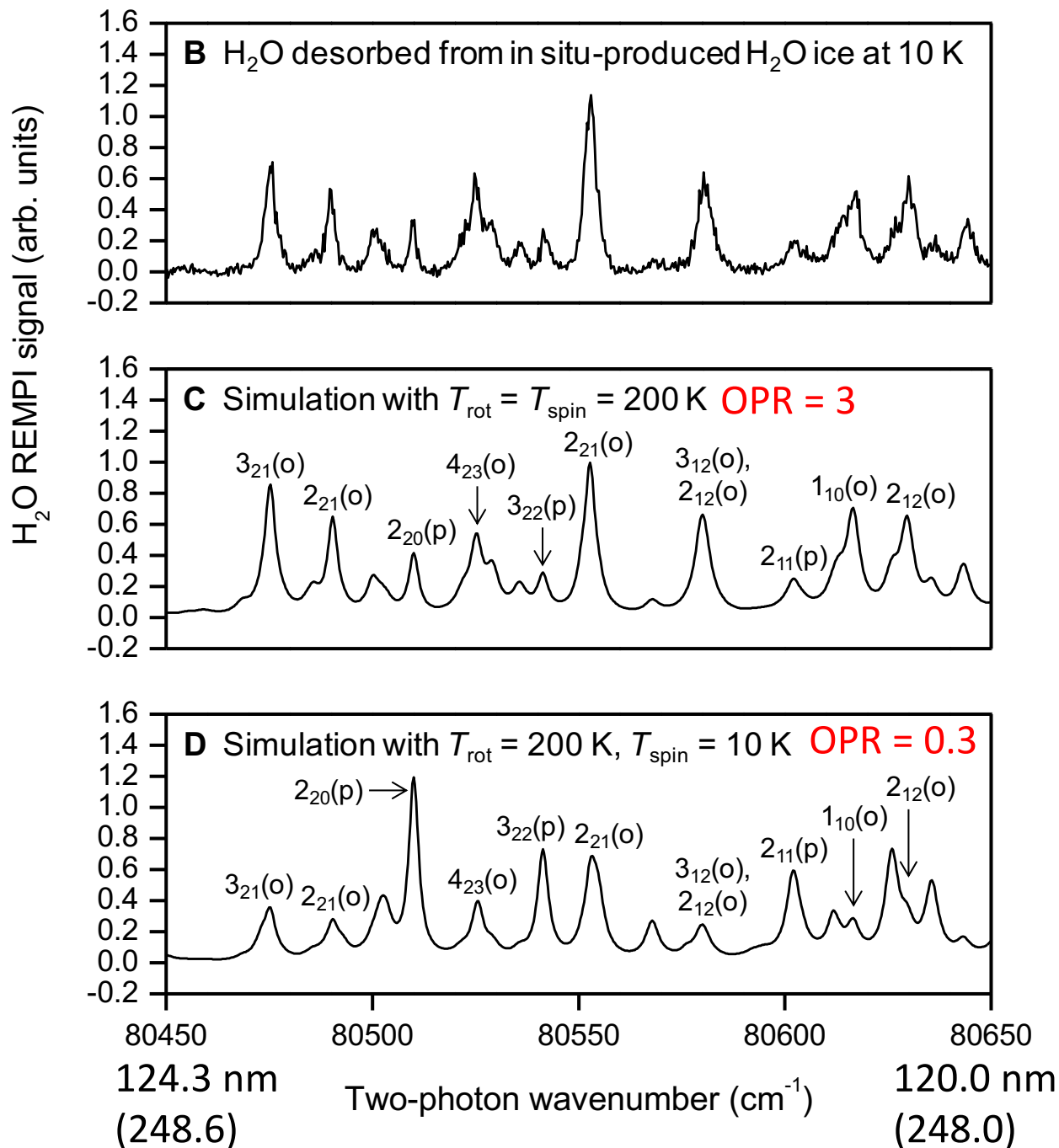
Photodesorption laser  
Unfocused (0.1 mJ pulse<sup>-1</sup>)  
10<sup>13</sup>-10<sup>14</sup> photons cm<sup>-2</sup> pulse<sup>-1</sup>

Amorphous H<sub>2</sub>O ice

Al substrate at 10 K



# REMPI rotational spectrum of H<sub>2</sub>O photodesorbed from ice at 10 K



Assignments:  $J_{KaKc}$  (ortho or para)

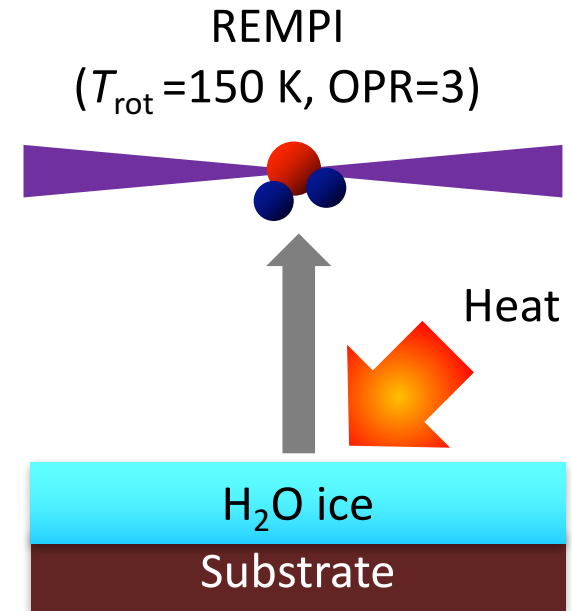
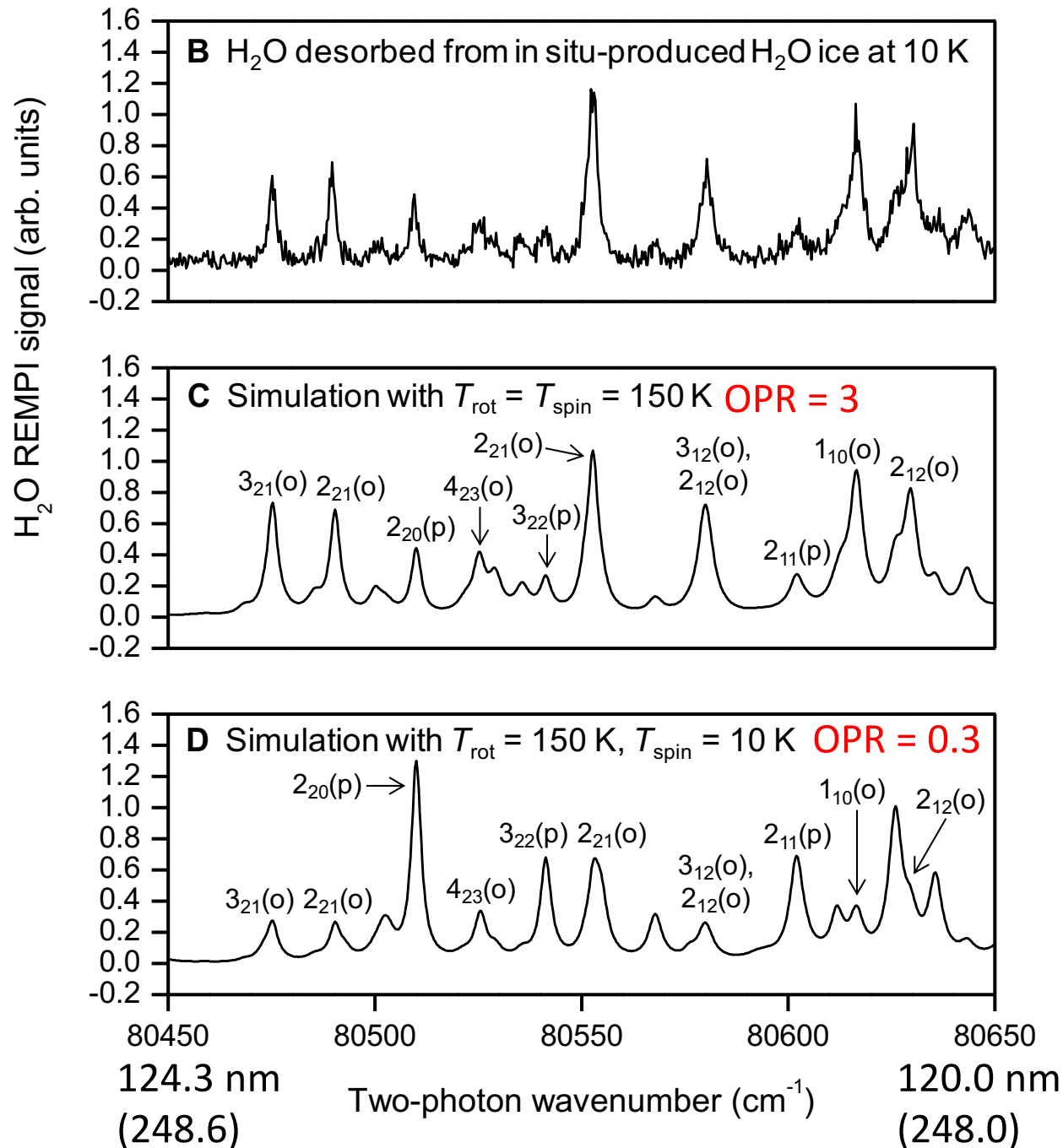
Stronger ortho-H<sub>2</sub>O lines than para-H<sub>2</sub>O lines.

Well-reproduced with OPR = 3.

↕  
Para-H<sub>2</sub>O signals are too strong with OPR = 0.3 (10 K).

**The OPR is not related to the ice formation temperature.**

# REMPI rotational spectrum of H<sub>2</sub>O thermally desorbed from ice at 150 K



Stronger ortho-H<sub>2</sub>O lines than para-H<sub>2</sub>O lines.

Well-reproduced with OPR = 3.

↕  
Para-H<sub>2</sub>O signals are too strong with OPR = 0.3 (10 K).

**The OPR is not related to the ice formation temperature.**

# Large difference of “rotational state of H<sub>2</sub>O” between the gas and solid phases.

The definition of spin temperature.

Spin and rotational degeneracy

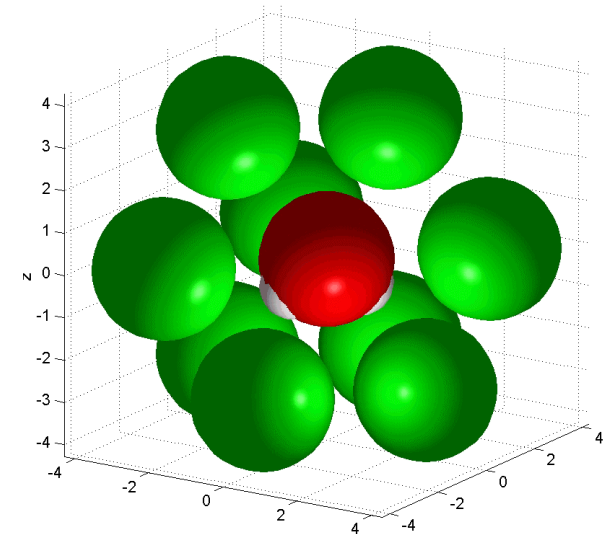
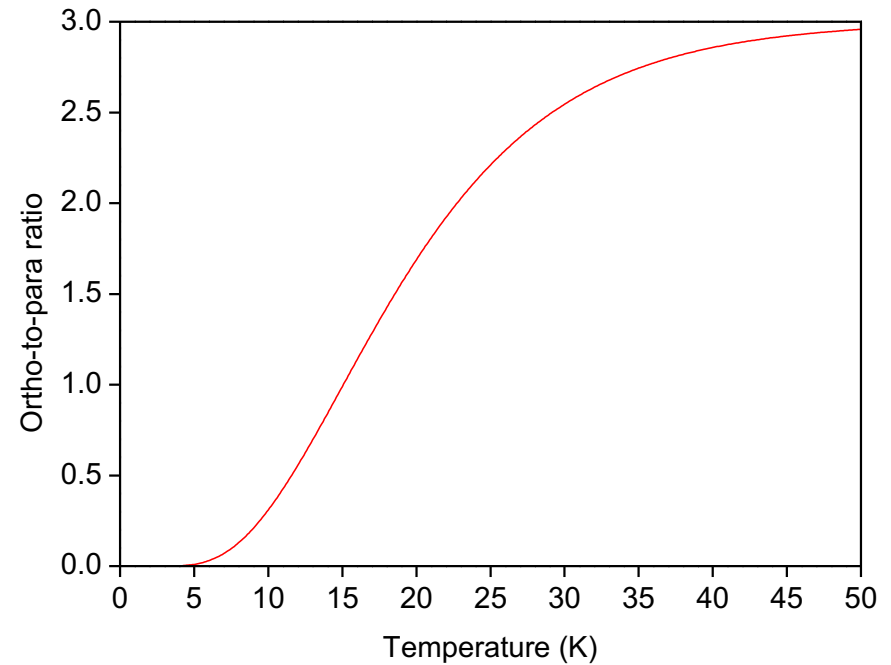
$$\text{OPR} = \frac{3 \sum (2J + 1) \exp \left[ \frac{-E_o(J_{K_a, K_c})}{k_B T_{\text{spin}}} \right]}{\sum (2J + 1) \exp \left[ \frac{-E_p(J_{K_a, K_c})}{k_B T_{\text{spin}}} \right]}$$

Rotational energy is  
important !!  
 $\Delta E = 34.2 \text{ K}$ .

In rotational-hindered system,  
the “rotational” energy difference ( $\Delta E_{\text{rot}}$ ) becomes small.

Even in a solid Ar matrix, the  $\Delta E_{\text{rot}}$  becomes small  $\sim 20 \text{ cm}^{-1}$ .

Turgeon et al., JPCA, 121, 1571 (2017).  
Michaut et al., Vib. Spec. 34 83 (2004).



# OPR- $T_{\text{spin}}$ curves in rotational hindered systems

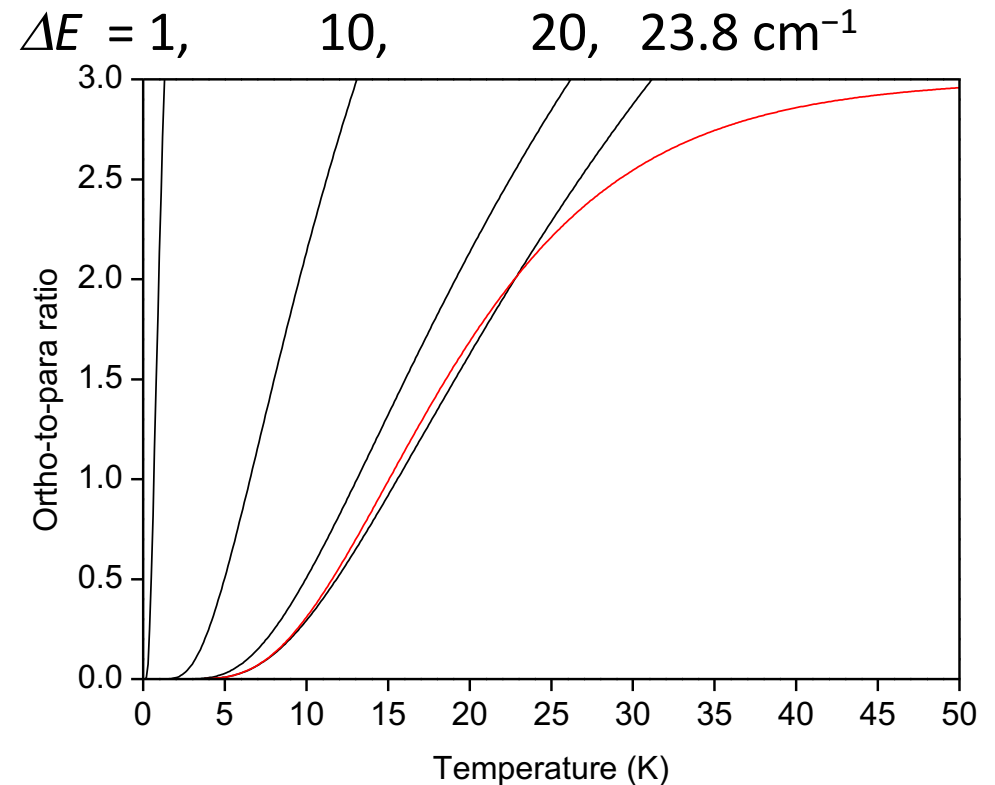
Approximated (crude) OPR- $T_{\text{spin}}$  curve;

$$\text{OPR} = \frac{3 \sum (2J + 1) \exp \left[ \frac{-E_o(J_{K_a, K_c})}{k_B T_{\text{spin}}} \right]}{\sum (2J + 1) \exp \left[ \frac{-E_p(J_{K_a, K_c})}{k_B T_{\text{spin}}} \right]}$$

$$\approx 9 \exp \left( \frac{-\Delta E}{k_B T_{\text{spin}}} \right)$$

Only consider the lowest two rotational levels.  
(since we cannot know all rotational levels.)

Rotationally hindered  $\longleftrightarrow$  Free rotation

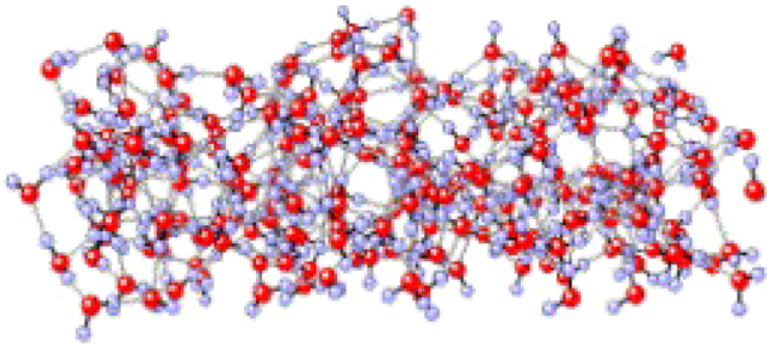


**As the  $\Delta E$  value decreased, the OPR reaches 3 at low temperature.  
→ The energies of ortho-, and para- $\text{H}_2\text{O}$  become comparable.**

**The  $T_{\text{spin}}$  curve only applies to  $\text{H}_2\text{O}$  in the gas phase (free rotors).**

In ice, H<sub>2</sub>O has a high barrier (6700 K) to rotation, because of hydrogen-bond network.

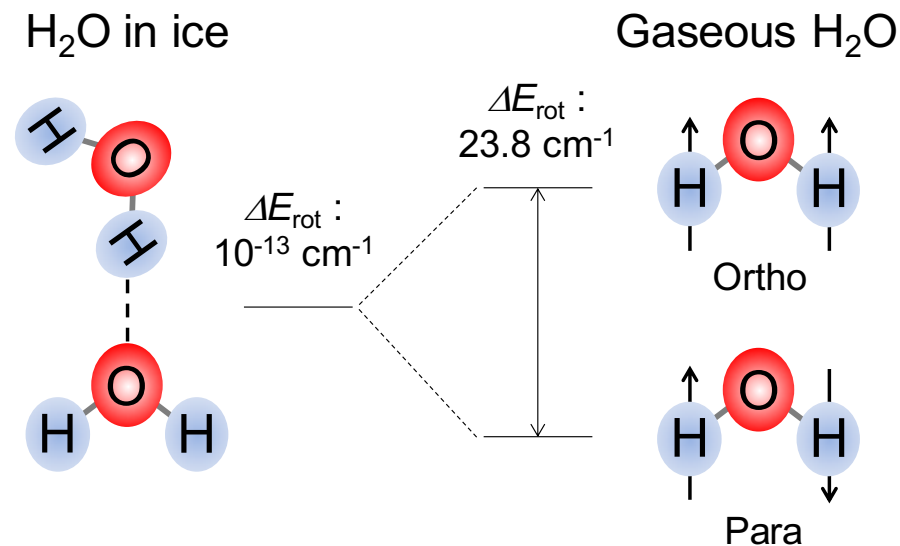
Wittebort et al., JACS. 110, 5668 (1988).



**Both ortho- and para-H<sub>2</sub>O are no longer free rotors.**  
**The quenched  $\Delta E$  value:  $10^{-13} \text{ cm}^{-1}$  ( $10^{-13} \text{ K}$ )**  
**(The same energy level !!)**

Buntkowsky et al., Z. Phys. Chem. 222, 1049 (2008).

**Ortho- and para-H<sub>2</sub>O are energetically comparable at 10 K in strongly rotational-hindered systems (ice).**



# Nuclear spin conversion of H<sub>2</sub>O

NSC: ortho-para state mixing by magnetic perturbation (e.g., paramagnetic O<sub>2</sub>).

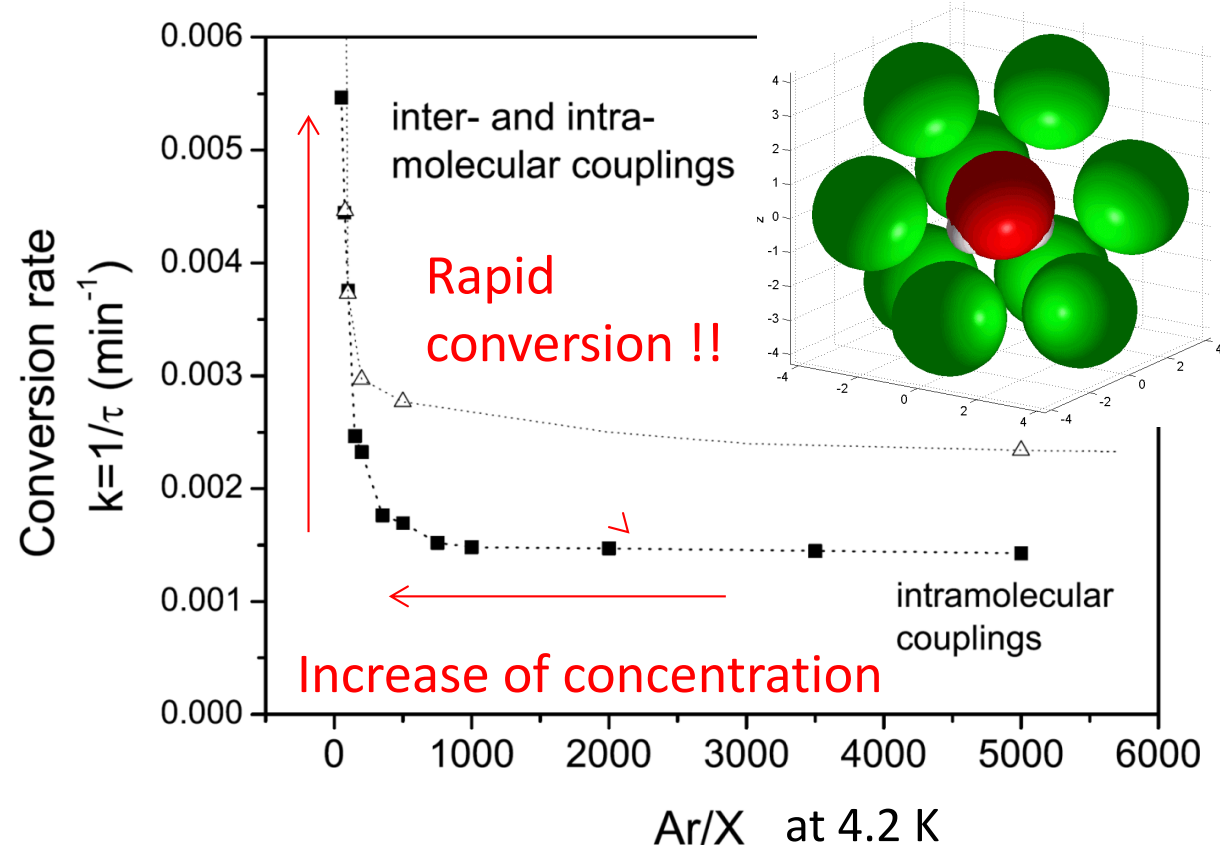
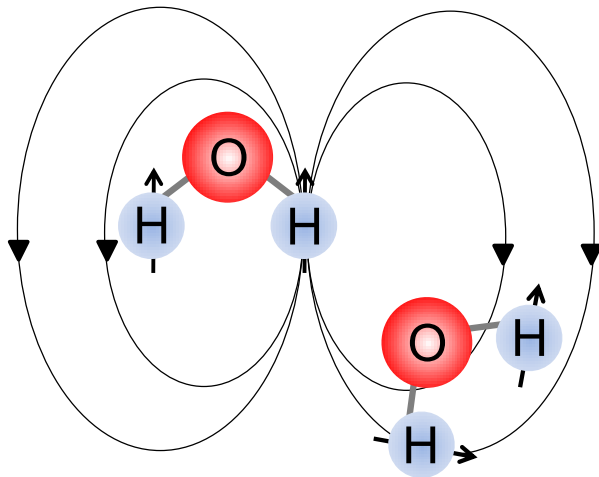
**Proton has a magnetic moment ( $10^{-26} \text{ J T}^{-1}$ )**

**Intermolecular proton-proton dipolar interaction can induce NSC.**

NSC of water (■) and methane (△) in an Ar matrix:

**Enhanced by intermolecular proton-proton dipolar interaction.**

$10^{-7}$  to  $10^{-6} \text{ cm}^{-1}$  ( $10^{-30} \text{ J}$ )  
(In terms of NMR, 10 kHz)



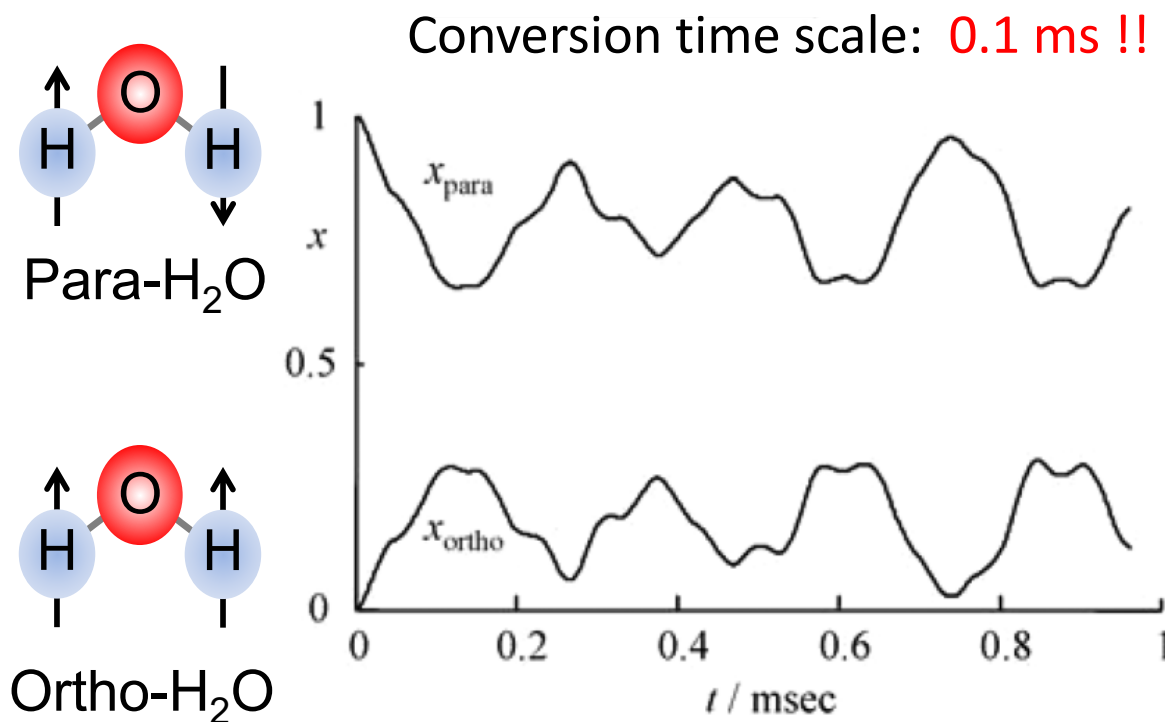
As a result of rotational quenching,  
fast nuclear spin conversion of H<sub>2</sub>O can occur in ice

**In ice, each proton feels magnetic fields created by all protons in ice.**



Ortho- and para-states are strongly mixed:

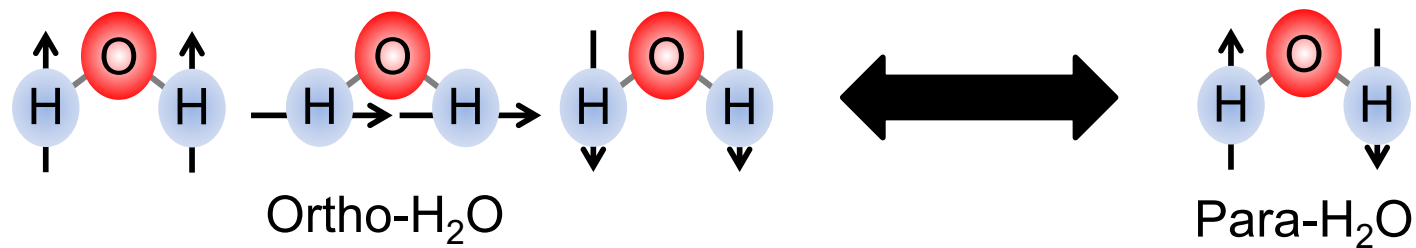
Intermolecular proton magnetic interaction  $\gg$  rotational energy difference ( $\Delta E_{\text{rot}}$ ).  
 $10^{-7}$  to  $10^{-6}$  cm<sup>-1</sup>  $\gg$   $10^{-13}$  cm<sup>-1</sup>.



Limbach et al. 2006,  
Chem. Phys. Chem., 7, 551.  
Buntkowsky et al. 2008,  
Z. Phys. Chem., 222, 1049.



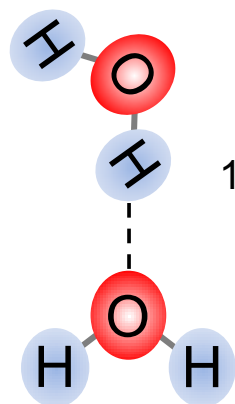
The OPR of H<sub>2</sub>O in ice is in dynamic equilibrium at the statistical value of 3.  
(The spin state of H<sub>2</sub>O is superposition of ortho- and para-H<sub>2</sub>O)



(1) Comparable energy of ortho- and para-H<sub>2</sub>O  
by rotational quenching

(2) Fast continuous  
ortho-para interconversion

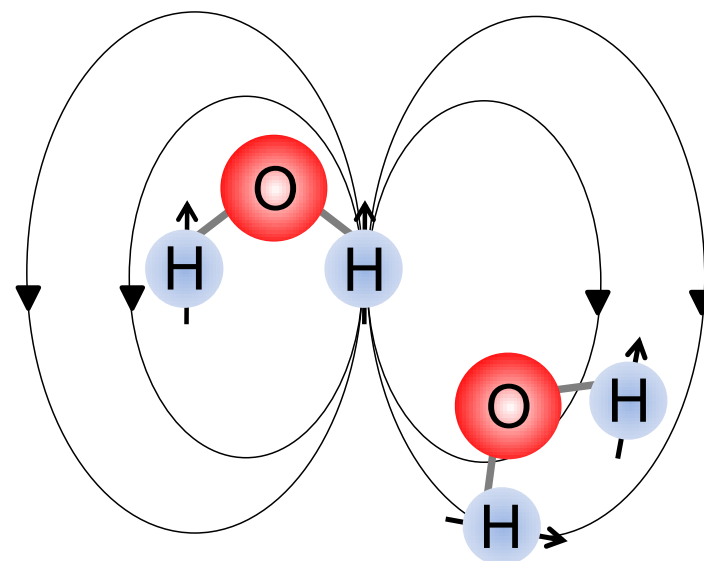
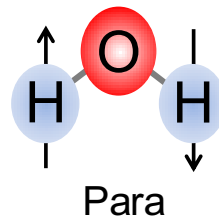
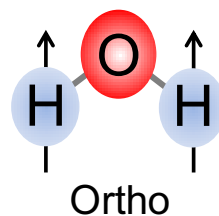
H<sub>2</sub>O in ice



$\Delta E_{\text{rot}} :$   
 $10^{-13} \text{ cm}^{-1}$

$\Delta E_{\text{rot}} :$   
 $23.8 \text{ cm}^{-1}$

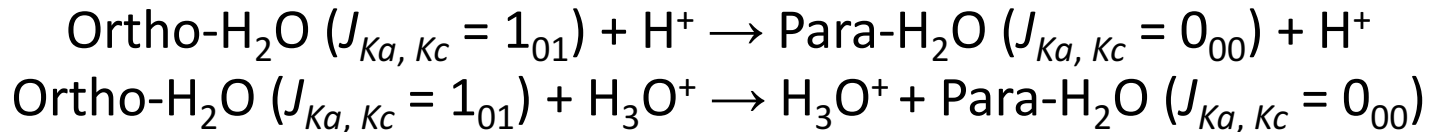
Gaseous H<sub>2</sub>O



The origin of the anomalous OPRs of interstellar H<sub>2</sub>O is still an open question.

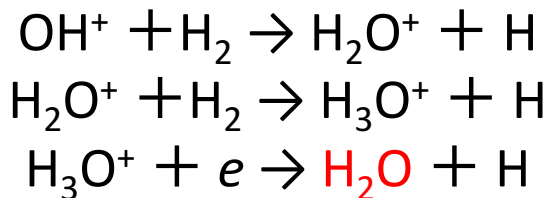
**Nuclear-spin effects in gas-chemistry are key.**

(1) NSC via ion-molecule reactions, e.g., H<sup>+</sup> or H<sub>3</sub>O<sup>+</sup>.



These two reactions can be endothermic at low temp. because  $\Delta E_{\text{rot}} = 34.2$  K.  
Might lead to para-enrichment of H<sub>2</sub>O.

(2) Gas-formation processes.



Since H<sub>2</sub> can be para-enriched in interstellar clouds,  
the H<sub>2</sub>O products might be also para-enriched.

H<sub>2</sub>O<sup>+</sup>, H<sub>3</sub>O<sup>+</sup> have also nuclear-spin isomers.

Nuclear-spin effects should be considered.

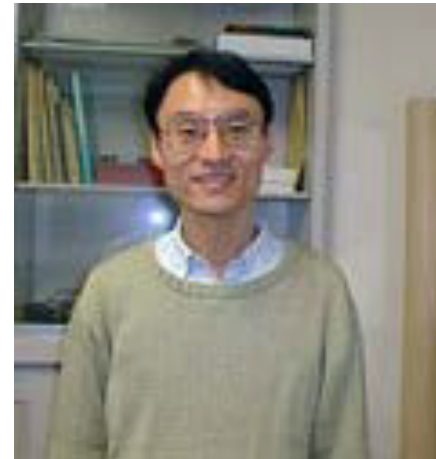
Accurate state-to-state rate coefficients at low temperature are unknown...

Gas-phase conversion: Lique et al., Int. Rev. Phys. Chem. 33, 125 (2014).  
Nuclear-spin selection rules: Gerlich et al., Proc. R. Soc. A. 364, 3007 (2006).

# Summary

- (1) H<sub>2</sub>O desorbed from ice at 10 K shows a statistical OPR of 3, even when the ice is produced in situ at 10 K.
- (2) Reinterpretation of previous observations is necessary.
- (3) Importance of the gas-phase chemistry.
  - (a) NSC via ion-molecule reactions, e.g., H<sup>+</sup> or H<sub>3</sub>O<sup>+</sup>.
  - (b) Gas-formation processes.

Hama, and Watanabe, Chem. Rev. 113, 8783 (2013).  
Hama, Kouchi, Watanabe, Science 351, 65-67 (2016).



Financial supports:  
JSPS (grant 24224012), MEXT (grant 25108002).



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日本学術振興会



# Another explanation for nuclear spin state of H<sub>2</sub>O in ice: NMR

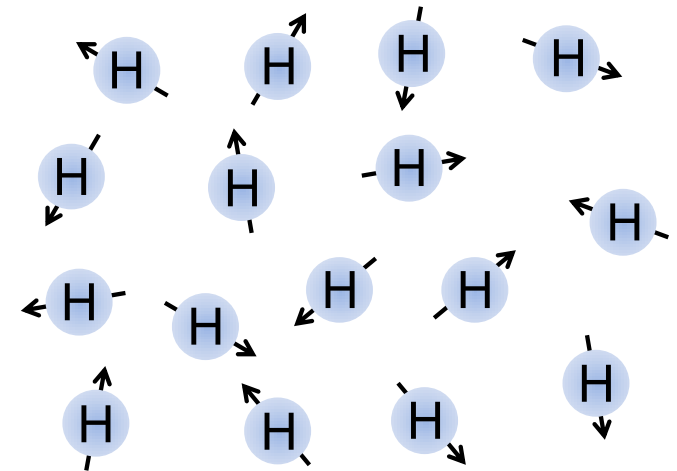
Water ice: a diamagnetic insulator (No macroscopic nuclear magnetization).



The direction of the proton spin angular momentum is distributed uniformly in ice.  
(No net magnetic moment).

Each proton feels  
magnetic fields created by all protons in ice, while  
Proton-proton dipolar interactions  $\ll$  thermal energy

$$10^{-30} \text{ J (10 kHz)} \ll 10^{-22} \text{ J (10 K)}$$



When an H<sub>2</sub>O molecule is desorbed from the ice into the gas phase,  
it is restored as “a coupled homonuclear spin-1/2 pair”, ortho ( $I=1$ ), or Para ( $I=0$ ).



Since the desorbed H<sub>2</sub>O molecules have sufficient energy,  
( $> 34.2 \text{ K}$ ,  $\Delta E_{\text{rot}}$  in the gas phase),  
the OPR can be statistical.

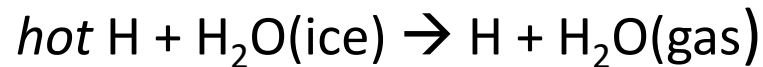
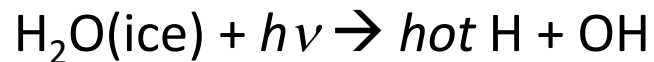
Detector

$T_{\text{rot}}$  and  $T_{\text{trans}}$  are different.

Not in thermodynamic equilibrium.

No local heating (non-equilibrium process)

“kick-out”



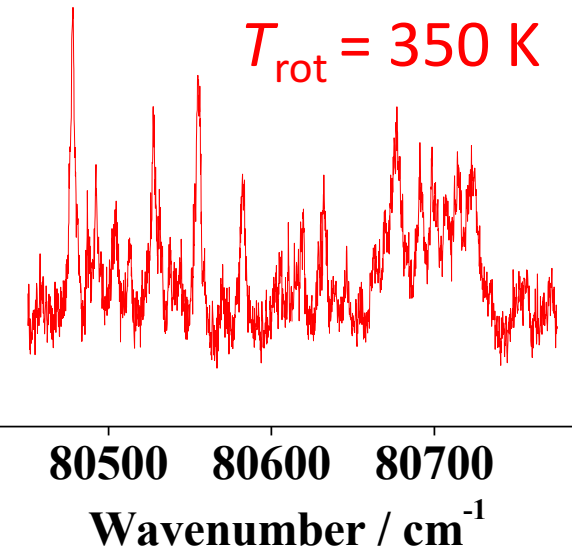
157 nm  
Photodissociation

2 mm

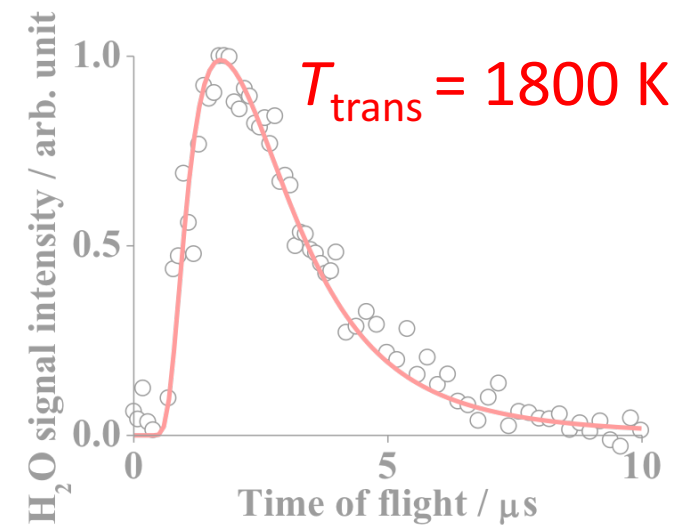
ice film

Au

### Rotational spectrum

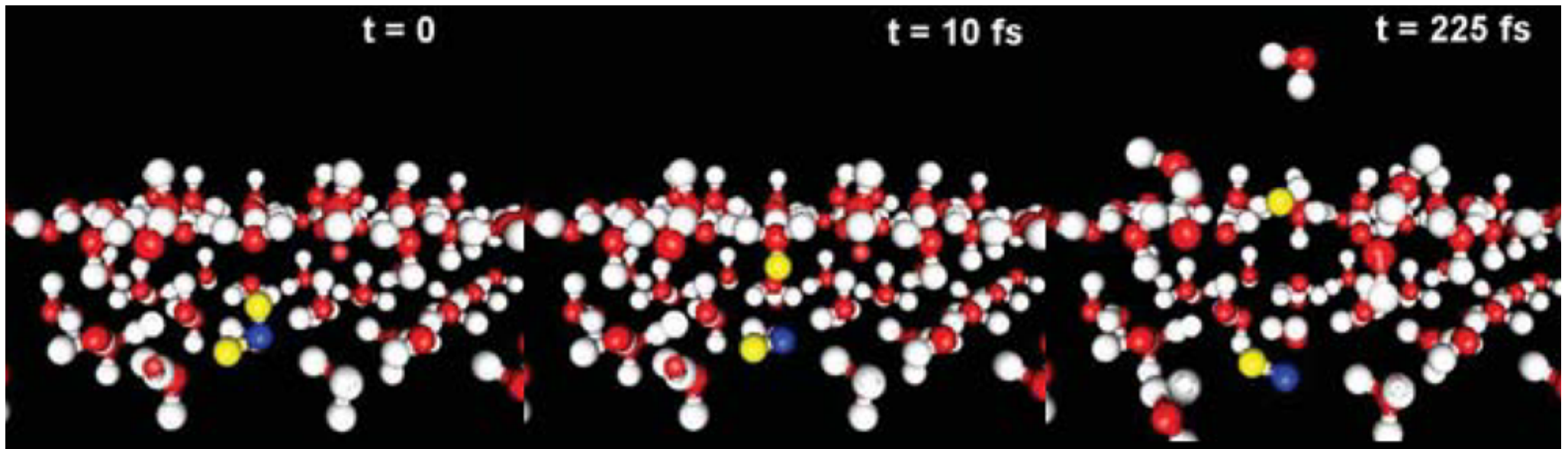
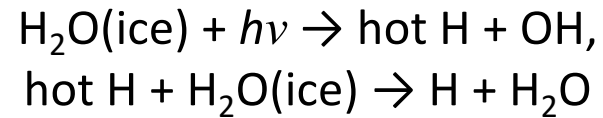


### Laser-delay spectrum



# Photodesorption mechanism

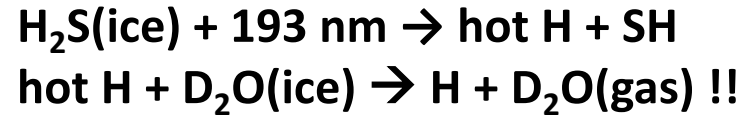
“kick-out” : An H<sub>2</sub>O molecule is desorbed without intramolecular bond dissociation by the momentum transfer from an energetic H atom photodissociated from a neighboring H<sub>2</sub>O.



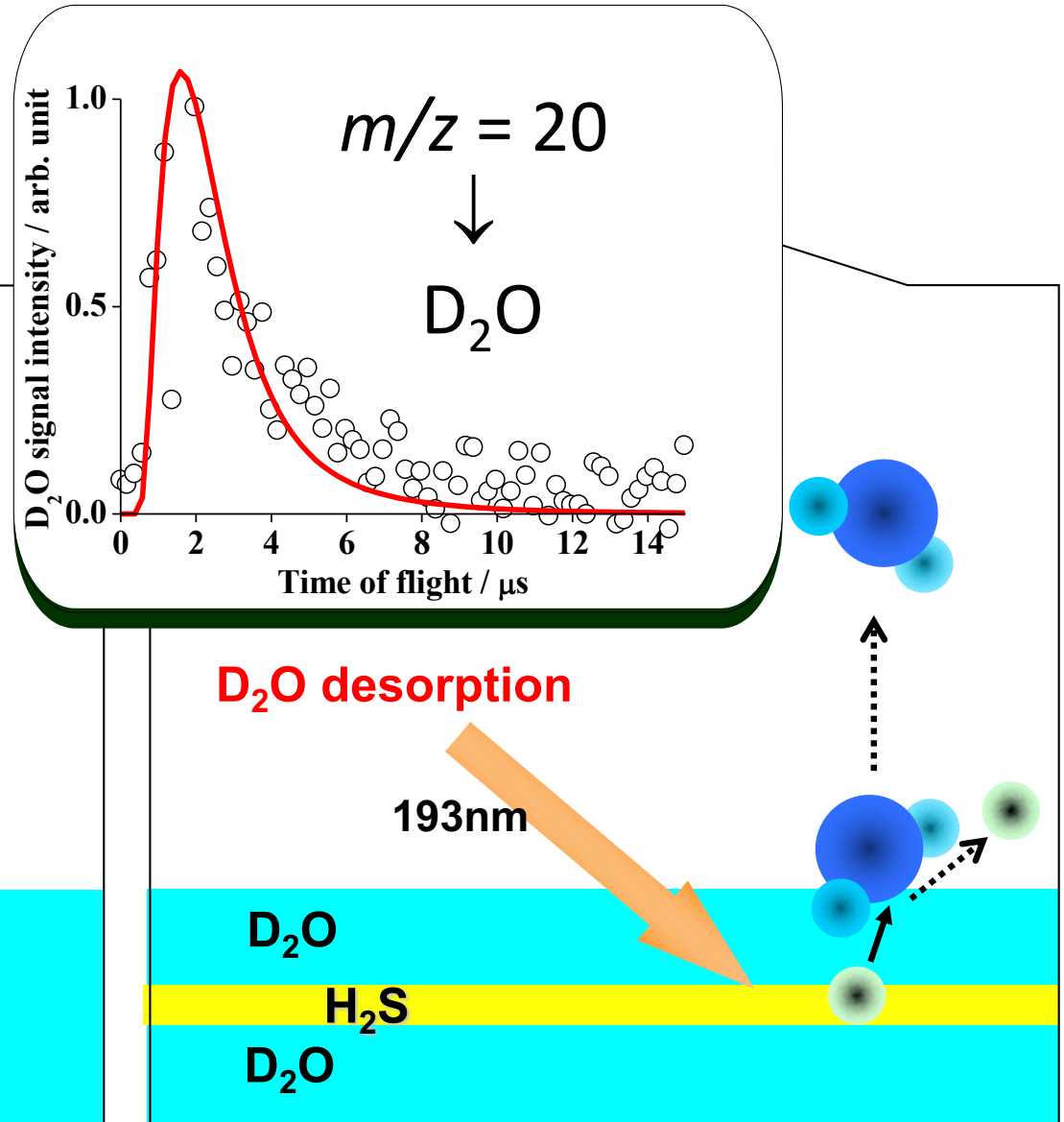
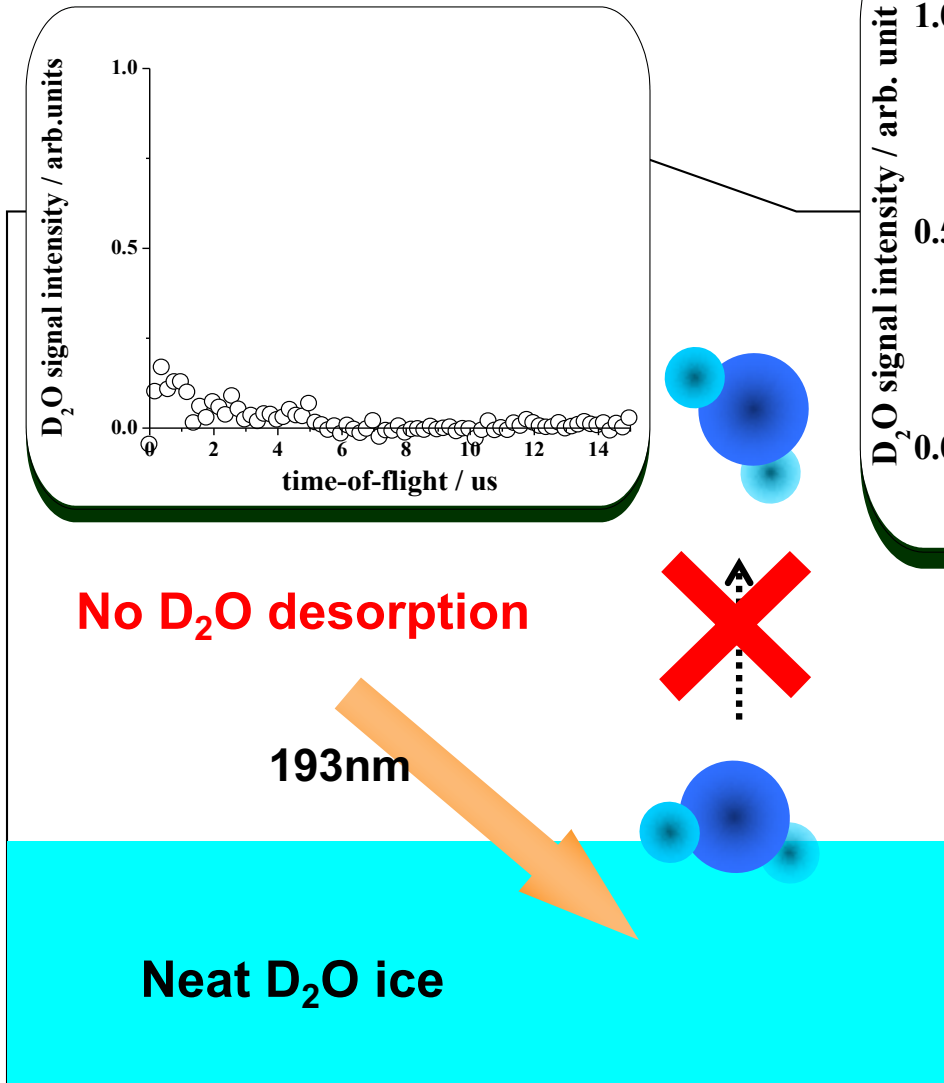
Andersson, and van Dishoeck, *Astron. Astrophys.* , 491, 907 (2008).



# kick-out: D<sub>2</sub>O desorption from “D<sub>2</sub>O + H<sub>2</sub>S layered ice” by 193 nm irradiation.



No D<sub>2</sub>O following 193 nm irradiation of pure D<sub>2</sub>O ice (no absorption).



# Resonance Enhanced Multi-Photon Ionization (REMPI) spectroscopy

Ordinary photoionization techniques cannot distinguish ortho/para  $\text{H}_2\text{O}$  ( $m/z = 18$ )

One photon ionization      Multiphoton ionization      2+1 REMPI

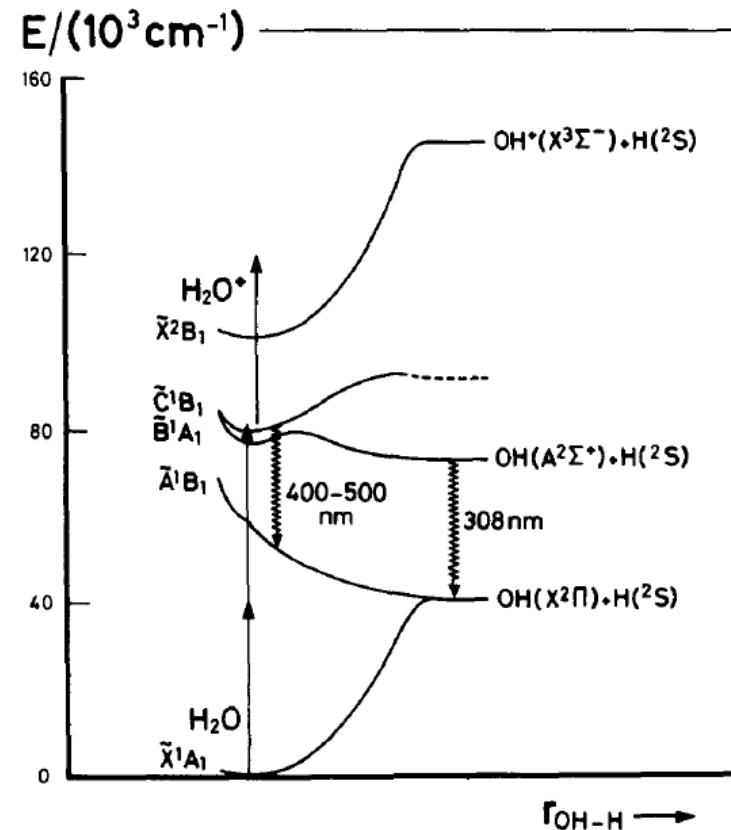
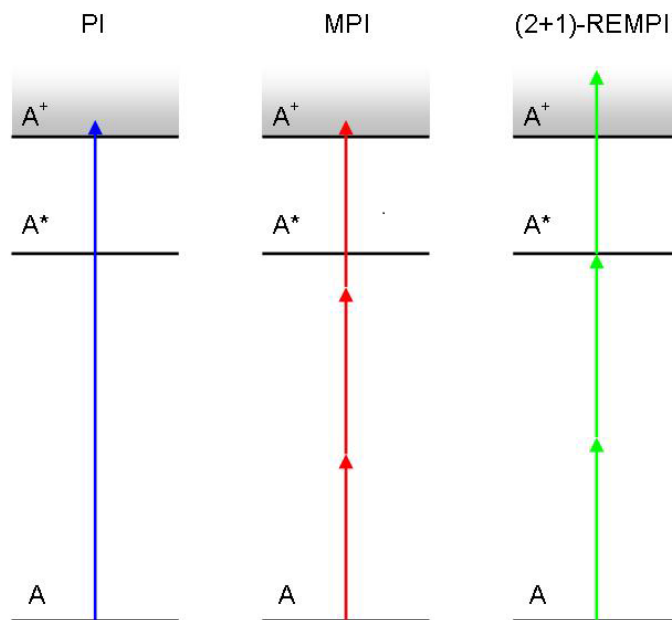


FIG. 1. Correlation diagram showing the processes involved (after Ref. 16).

In 2+1 REMPI of  $\text{H}_2\text{O}$ , 2-photon absorption excites population from the ground  $X^1A_1$  ( $v=0, J_{Ka,Kc}$ ) state to the intermediate  $C^1B_1$  ( $v'=0, J'_{Ka',Kc'}$ ) state. The absorption of one further photon transfers population into the ionization continuum.

Since the REMPI transition is rotationally (i.e., OPR) resolved, we can get the rotational and spin temperature of desorbed  $\text{H}_2\text{O}$ .

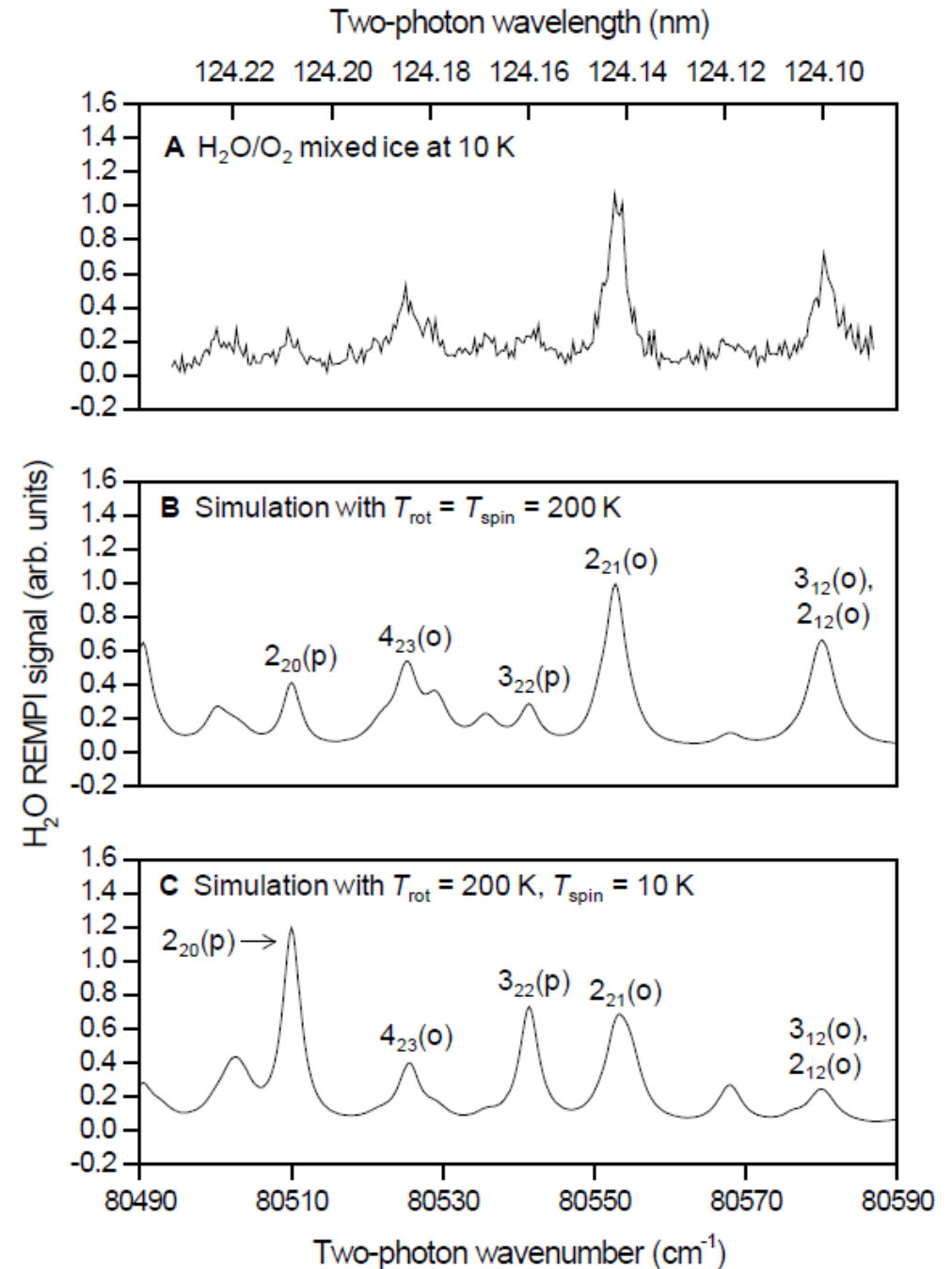
# Paramagnetic catalytic effects on the NSC of H<sub>2</sub>O during photodesorption.

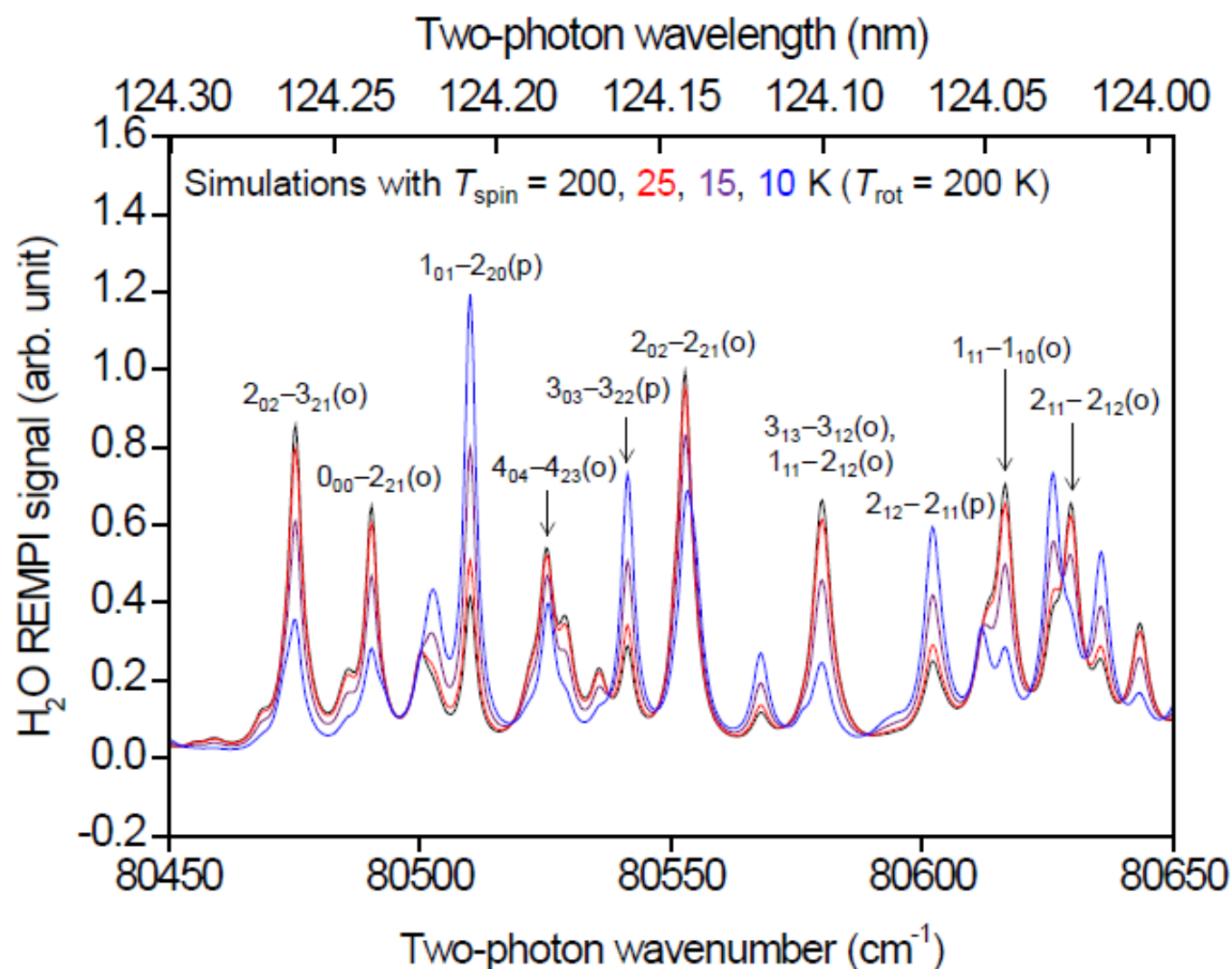
REMPI spectra of photodesorbed H<sub>2</sub>O from H<sub>2</sub>O/O<sub>2</sub> (1:1) mixed ice at 10 K.

Stronger ortho-H<sub>2</sub>O lines than para-H<sub>2</sub>O lines.

Reproduced by the simulation with  $T_{rot} = T_{spin} = 200$  K.

Although the electron magnetic moment is about -658 times larger than the proton magnetic moment, the timescale for photodesorption (femto-s) is much shorter than that for NSC through magnetic interactions (micro-sec).

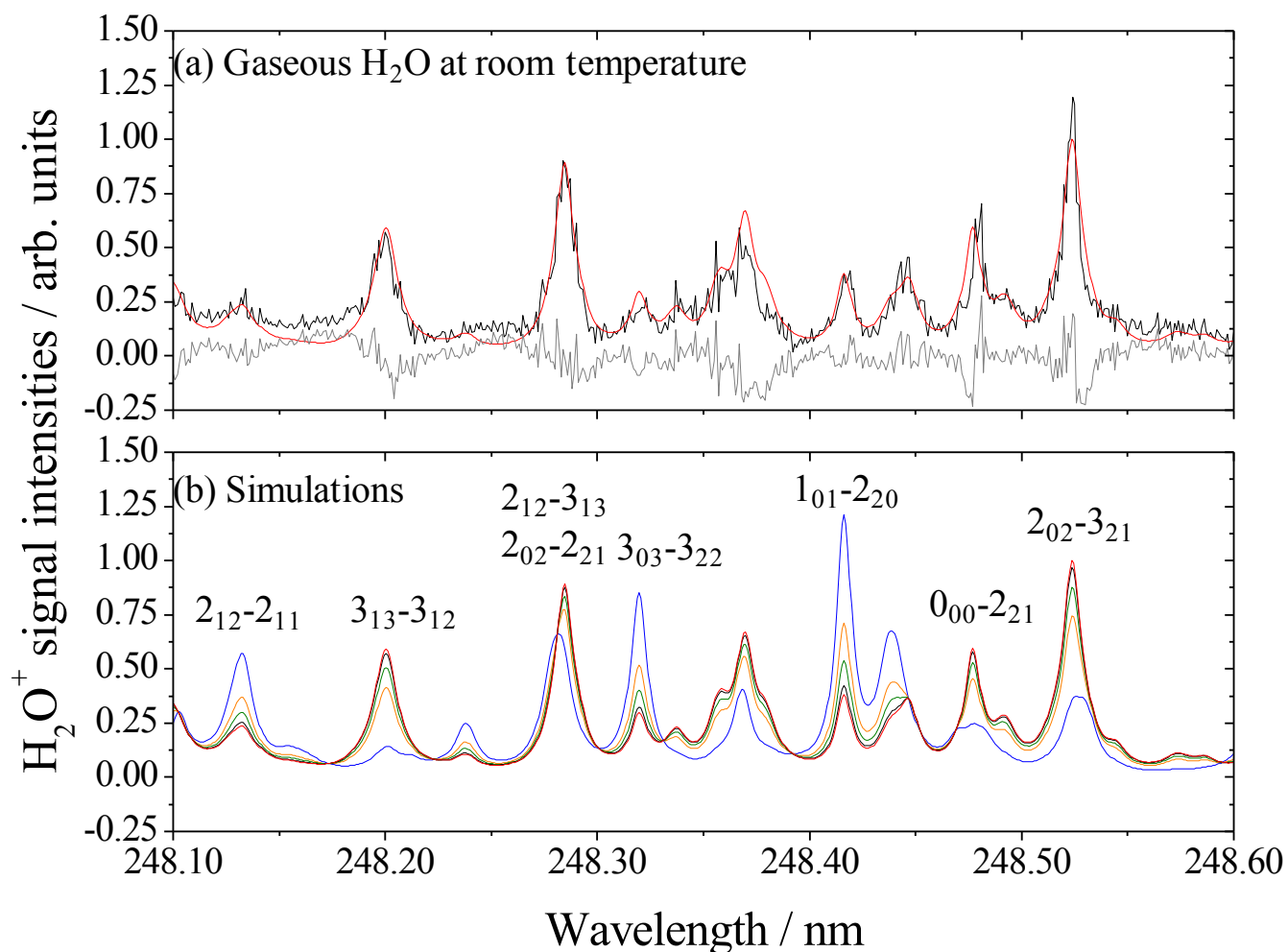




**Fig. S3: REMPI spectral simulations.**  $T_{\text{spin}} = 200$  (black), 25 (red), 15 (purple), and 10 (blue) K.  $T_{\text{rot}}$  is fixed at 200 K.  $T_{\text{spin}}$  and  $T_{\text{rot}}$  represent nuclear spin and rotational temperatures, respectively. Indications ( $J'_{Ka',Kc'} - J_{Ka,Kc}$ ) are rotational assignments of the  $\tilde{C}^1B_1(v=0) - \tilde{X}^1A_1(v=0)$  transition in  $\text{H}_2\text{O}$ , where “o” and “p” denote ortho and para, respectively.

2+1 REMPI spectrum of H<sub>2</sub>O via the  $C^1B_1(v=0) \leftarrow X^1A_1(v=0)$  transition.

$(J'_{Ka',Kc'} \leftarrow J_{Ka,Kc})$  are rotational assignments.



$J_{KaKc}$  (ortho or para)

$T_{\text{rot}} = T_{\text{spin}} = 300 \text{ K}$

Gray; difference

$T_{\text{rot}} = T_{\text{spin}} = 300 \text{ K}$

$T_{\text{rot}} = 300 \text{ K}, T_{\text{spin}} = 30 \text{ K}$

$T_{\text{rot}} = 300 \text{ K}, T_{\text{spin}} = 20 \text{ K}$

$T_{\text{rot}} = 300 \text{ K}, T_{\text{spin}} = 15 \text{ K}$

$T_{\text{rot}} = 300 \text{ K}, T_{\text{spin}} = 8 \text{ K}$